

For a multi-variables function:

$$f(x, y) = f(x_0, y_0) + \frac{(\partial f / \partial x)^{(1)}}{1!} (x - x_0) + \dots + \frac{(\partial f / \partial y)^{(1)}}{1!} (y - y_0) + \dots \quad (10)$$

And,

$$\vec{B}[B_x(x, y, z), B_y(x, y, z), B_z(x, y, z)] = \vec{i}B_x + \vec{j}B_y + \vec{k}B_z \quad (11)$$

Using Taylor Formula,

$$B_x(x, y, z) = B_x(x_0, y_0, z_0) + \frac{\partial B_x}{\partial x} (x - x_0) + \dots + \frac{\partial B_x}{\partial y} (y - y_0) + \dots + \frac{\partial B_x}{\partial z} (z - z_0) + \dots \quad (12)$$

And for  $\vec{r}(\nabla \cdot \vec{B})$ ,

$$\vec{r}(\nabla \cdot \vec{B}) = (\vec{i}x + \vec{j}y + \vec{k}z) \cdot \left( \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} + \frac{\partial B_z}{\partial z} \right), \quad (13)$$

and,

$$(\vec{r} \cdot \nabla) \vec{B} = \left( x \frac{\partial}{\partial x} + y \frac{\partial}{\partial y} + z \frac{\partial}{\partial z} \right) \cdot (\vec{i}B_x + \vec{j}B_y + \vec{k}B_z). \quad (14)$$

Which one meets the demand of equation ?? is extremely obvious. The proof can always be used in the first order approximation. Like when we try to calculate the magnetic drift in plasma physics.

**Proof of  $\nabla \times (\nabla \times \vec{A}) = \nabla(\nabla \cdot \vec{A}) - (\nabla \cdot \nabla) \vec{A}$ .** First,

$$\nabla \times \vec{B} = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ B_x & B_y & B_z \end{vmatrix} = \left( \frac{\partial B_z}{\partial y} - \frac{\partial B_y}{\partial z} \right) \vec{i} + \left( \frac{\partial B_x}{\partial z} - \frac{\partial B_z}{\partial x} \right) \vec{j} + \left( \frac{\partial B_y}{\partial x} - \frac{\partial B_x}{\partial y} \right) \vec{k} \quad (15)$$

$$\nabla \times (\nabla \times \vec{A}) = \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ \frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} & \frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} & \frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \end{vmatrix}. \quad (16)$$

$$\nabla(\nabla \cdot \vec{A}) = \left( \vec{i} \frac{\partial}{\partial x} + \vec{j} \frac{\partial}{\partial y} + \vec{k} \frac{\partial}{\partial z} \right) \left( \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z} \right), \quad (17)$$

$$(\nabla \cdot \nabla) \vec{A} = \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) (\vec{i}A_x + \vec{j}A_y + \vec{k}A_z). \quad (18)$$

Then, it is quite easy to use preliminary school calculating method to see their mathematical relation. **All proofs above can be better written in a tensor form.**

$$\nabla \cdot \mathbf{A} = \sum_i \frac{\partial A_i}{\partial x_i} = \partial_i A^i. \quad (19)$$

$$\begin{aligned} \nabla \times \mathbf{A} &= \sum_i \left( \frac{\partial A_k}{\partial x_j} - \frac{\partial A_j}{\partial x_k} \right) \hat{e}_i \\ &= \sum_{ijk} \varepsilon_{ijk} \frac{\partial A_j}{\partial x_i} \hat{e}_k \\ &= \varepsilon^{ijk} \partial_i A_j = \varepsilon^{kij} \partial_i A_j. \end{aligned} \quad (20)$$

# Additional Vector Analysis

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**Proof 1 of dyadic vector.**

$$\begin{aligned} \vec{a} \times (\vec{b} \cdot \vec{c}) \vec{d} &= \vec{a} \times \vec{d} (\vec{b} \cdot \vec{c}) \\ &= \vec{e} (\vec{b} \cdot \vec{c}) \\ &= (\vec{e}_2 - \vec{e}_1) [(\vec{b}_2 - \vec{b}_1) \cdot \vec{c}] \\ &= [\vec{e}_2 \vec{b}_2 - \vec{e}_2 \vec{b}_1 - \vec{e}_1 \vec{b}_2 + \vec{e}_1 \vec{b}_1] \cdot \vec{c}. \end{aligned} \quad (1)$$

**Proof of  $\vec{k} = -(\vec{b} \cdot \nabla) \vec{b}$ .** The basic:

$$\begin{aligned} \vec{b} &= \vec{i}b_x + \vec{j}b_y + \vec{k}b_z, \\ \nabla &= \vec{i} \frac{\partial}{\partial x} + \vec{j} \frac{\partial}{\partial y} + \vec{k} \frac{\partial}{\partial z}. \end{aligned} \quad (2)$$

For  $\vec{b} \cdot \nabla$  as a united operator:

$$\vec{b} \cdot \nabla = b_x \frac{\partial}{\partial x} + b_y \frac{\partial}{\partial y} + b_z \frac{\partial}{\partial z}. \quad (3)$$

The curvature  $\vec{k}$  (not the same  $\vec{k}$  as the unit vector) can be defined as  $d\vec{b}/dl$  (which is very straightforward to understand):

$$\frac{d\vec{b}}{dl} = \frac{\vec{i}db_x + \vec{j}db_y + \vec{k}db_z}{dl}. \quad (4)$$

Multiple the operator  $\vec{b} \cdot \nabla$  and  $\vec{b}$ , focus on the first term (since the second and the third term use the same approach):

$$\vec{i} \left( b_x \frac{\partial b_x}{\partial x} + b_y \frac{\partial b_x}{\partial y} + b_z \frac{\partial b_x}{\partial z} \right). \quad (5)$$

We have:

$$\begin{aligned} db_x &= \frac{\partial b_x}{\partial x} dx + \frac{\partial b_x}{\partial y} dy + \frac{\partial b_x}{\partial z} dz, \\ \frac{db_x}{dl} &= \frac{\partial b_x}{\partial x} \frac{dx}{dl} + \frac{\partial b_x}{\partial y} \frac{dy}{dl} + \frac{\partial b_x}{\partial z} \frac{dz}{dl}. \end{aligned} \quad (6)$$

Using the definition of the unit vector  $\vec{b}$  and a very easy triangle relation, we have:

$$\frac{dx}{dl} = b_x = \frac{b_x}{|b|}, \quad (7)$$

combine equation ?? and equation ??, it is very easy to get:

$$b_x \frac{\partial b_x}{\partial x} + b_y \frac{\partial b_x}{\partial y} + b_z \frac{\partial b_x}{\partial z}, \quad (8)$$

which happens to be the norm of the curvature vector  $\vec{k}$  in the  $\vec{i}$  direction, which is shown in equation ?. This is also why the equation should be written at the form of  $\vec{k} = -(\vec{b} \cdot \nabla) \vec{b}$ , instead of  $\vec{k} = -(\vec{b} \cdot \nabla) \vec{b}$  or something else. This will soon be proven next.

**Proof of why the Taylor Formula of a vector should be written at the form of  $(\vec{r} \cdot \nabla) \vec{B}_0$  instead of  $\vec{r}(\nabla \cdot \vec{B}_0)$ .**

The ordinary Taylor Formula is written as:

$$f(x) = \frac{f(x_0)}{0!} + \frac{f'(x_0)}{1!} (x - x_0) + \dots + \frac{f^{(n)}(x_0)}{n!} (x - x_0)^n + \dots \quad (9)$$